

MHD Mixed Convection of Viscoelastic Nanofluid Flow due to Constant Heat Flux

Rahimah Mahat^{1,*}, Muhammad Saqib², Imran Ulah³, Sharidan Shafie², Sharena Mohamad Isa⁴

- ¹ Advanced Facilities Engineering Technology Research Cluster (AFET), Universiti Kuala Lumpur Malaysian Institute of Industrial Technology, Persiaran Sinaran Ilmu, Bandar Seri Alam, 81750 Masai, Johor, Malaysia
² Department of Mathematical Sciences, Faculty of Science, Universiti Teknologi Malaysia, 81310 Johor Bahru, Malaysia
³ College of Civil Engineering, National University of Sciences and Technology, 44000, Islamabad, Pakistan
⁴ Manufacturing Engineering Technology Section, Universiti Kuala Lumpur Malaysia Italy Design Institute (UniKL MIDI), Kuala Lumpur, 56100, Malaysia

ABSTRACT

A mathematical model for magnetohydrodynamics (MHD) viscoelastic nanofluid flow approaching a linear horizontal circular cylinder has been constructed. The cylinder with constant heat flux is illustrated in the analysis with magnetic field. We used the Tiwari and Das Nanofluid model in this analysis to learn more about nanofluid effects, where sodium carboxymethyl cellulose with nanoparticles of copper (Cu) was selected as a based fluid. Suitable transformations are used to turn dimensional linear equations into dimensionless expressions. To address the governing dimensionless issues, the Keller box method approach is implemented and executed. The effects of a few chosen factors on flow and heat transmission are investigated. The coefficients of skin friction and heat transmission are provided and analysed. The acquired findings are compared to the existing data in the limiting scenario, and there is a lot of consistency. It was found that velocity, temperature, skin friction and heat transfer coefficients of viscoelastic nanofluid depend strongly on viscosity and thermal conductivity together with magnetic field.

Keywords:

Magnetohydrodynamics (MHD);
Viscoelastic; Nanofluid; Constant heat flux

Received: 18 January 2022

Revised: 26 March 2022

Accepted: 28 April 2022

Published: 25 May 2022

1. Introduction

Heat transfer enhancement, augmentation, or intensification is a term used by researchers to describe a variety of strategies for improving heat transfer performance. Therefore, the addition of nanofluid in fluid flow is to enhance the heat transfer. The study on nanofluids have been studied widely by Waini *et al.*, Mahat *et al.*, Patil *et al.*, Murad *et al.*, and Shorbagy *et al.* [1-7]. The study of the flow of an electrically conducting fluid in a magnetic field is known as magnetohydrodynamics (MHD). Fundamentally, magnetic fields produce currents in a flowing conductive fluid, which tend to form Lorentz force drag forces on the fluid flow. The fluid flow will be affected by the Lorentz force. MHD is being used in astronomy and geophysics, nuclear fission and fusion, metallurgy, and direct energy conversion, among other fields. Many researchers, such as Jamel *et al.*, Hussain *et al.*, Tlili *et al.*, Wakif *et al.*, Hatami *et al.*, Cao *et al.*, Liaqat *et al.*, El-Shorbagy *et al.* and Jamel *et al.* [8-14], have undertaken extensive research on MHD of nanofluids. The benefits of MHD's influence on mixed

* Corresponding author.

E-mail address: rahimahm@unikl.edu.my

convection flow via a horizontal circular cylinder are evident in the technology's growth in industrial and technical applications. In the current research, the effects of MHD mixed convection of viscoelastic nanofluid past a horizontal circular cylinder is theoretically investigated. The effects of magnetic field, nanoparticle volume fraction and how these affect the thermal characteristics of the system are of particular interest.

2. Methodology

The Cartesian coordinate (x, y) is used and the dimensional gravitational acceleration is $g_x = g \sin(\bar{x}/a)$, where \bar{x} is the distance from the lower stagnation point. The dimensional velocity outside the boundary layer is $\bar{u}_e(\bar{x}) = U_\infty \sin(\bar{x}/a)$ and the constant free stream velocity is $(1/2)U_\infty$ as mentioned by Merkin [15]. Tiwari and Das model [16] has been chosen in this study and the model is defined as a single-phase model that use Brickman viscosity model. Under the above assumptions and by considering the nanofluid model, the dimensional governing equations of momentum equation and energy equation can be expressed as

$$\frac{\partial \bar{u}}{\partial \bar{x}} + \frac{\partial \bar{v}}{\partial \bar{y}} = 0, \tag{1}$$

$$\bar{u} \frac{\partial \bar{u}}{\partial \bar{x}} + \bar{v} \frac{\partial \bar{u}}{\partial \bar{y}} = \bar{u}_e \frac{\partial \bar{u}_e}{\partial \bar{x}} + \frac{\mu_{nf}}{\rho_{nf}} \frac{\partial^2 \bar{u}}{\partial \bar{y}^2} - \frac{k_o}{\rho_{nf}} \left[\frac{\partial}{\partial \bar{x}} \left(\bar{u} \frac{\partial^2 \bar{u}}{\partial \bar{y}^2} \right) + \bar{v} \frac{\partial^3 \bar{u}}{\partial \bar{y}^3} - \frac{\partial \bar{u}}{\partial \bar{y}} \frac{\partial^2 \bar{u}}{\partial \bar{x} \partial \bar{y}} \right] + g \beta_{nf} (T - T_\infty) \sin\left(\frac{\bar{x}}{a}\right) - \sigma B_0^2 \bar{u}, \tag{2}$$

$$\bar{u} \frac{\partial T}{\partial \bar{x}} + \bar{v} \frac{\partial T}{\partial \bar{y}} = \frac{k}{\rho C_p} \frac{\partial^2 T}{\partial \bar{y}^2}, \tag{3}$$

with the boundary conditions

$$\begin{aligned} \bar{u} = 0, \quad \bar{v} = 0, \quad T = -\frac{q_w}{k_{nf}} \quad & \text{at} \quad \bar{y} = 0, \quad \bar{x} \geq 0, \\ \bar{u} = \bar{u}_e(\bar{x}), \quad \frac{\partial \bar{u}}{\partial \bar{y}} = 0, \quad T = T_\infty \quad & \text{at} \quad \bar{y} \rightarrow \infty, \quad \bar{x} \geq 0, \end{aligned} \tag{4}$$

where k_{nf} is the thermal conductivity of nanofluid, T is the fluid temperature and q_w is the constant heat flux. The numerical values of the thermophysical properties of base fluid and nanoparticles are given in Table 1.

Table 1
 Thermophysical properties of nanoparticles and base fluid

Physical Properties	$\rho(\text{kg } m^{-3})$	$C_p(\text{J kg}^{-1}\text{K}^{-1})$	$k(\text{Wm}^{-1}\text{K}^{-1})$	$\beta \times 10^5 (\text{K}^{-1})$
Base Fluid(CMC)	997.1	4179	0.613	21
Nanoparticle (Cu)	8933	385	401	1.67

The dimensionless variables are introduced to simplify the complexity of the governing equations. Based on Anwar *et al.* [17], the dimensionless variables are defined as

$$\begin{aligned}
 x &= \bar{x}/a, \quad y = Re^{1/2}(\bar{y}/a), \quad u = \bar{u}/U_\infty, \quad v = Re^{1/2}(\bar{v}/U_\infty), \\
 u_e(x) &= \bar{u}_e(\bar{x})/U_\infty, \quad \theta = (T - T_\infty)/(T_f - T_\infty),
 \end{aligned} \tag{5}$$

where Re is Reynolds number. By substituting Eq. (5) into Eqs (1) - (3), the dimensionless system below is yielded

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \tag{6}$$

$$\left[(1-\phi) + \phi \frac{\rho_s}{\rho_f} \right] \left[u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right] = \left[(1-\phi) + \phi \frac{\rho_s}{\rho_f} \right] \sin x \cos x + \frac{1}{(1+\phi)^{2.5}} \frac{\partial^2 u}{\partial y^2} \tag{7}$$

$$-K \left[\frac{\partial}{\partial x} \left(u \frac{\partial^2 u}{\partial y^2} \right) + v \frac{\partial^3 u}{\partial y^3} - \frac{\partial u}{\partial y} \frac{\partial^2 u}{\partial x \partial y} \right] + \left[(1-\phi) + \phi \frac{(\rho\beta)_s}{(\rho\beta)_f} \right] \lambda \theta \sin(x) - \frac{\sigma B_0^2 a}{\rho_f U_\infty} u,$$

$$\left[(1-\phi) + \phi \frac{(\rho C_p)_s}{(\rho C_p)_f} \right] \left[u \frac{\partial \theta}{\partial x} + v \frac{\partial \theta}{\partial y} \right] = \frac{(k_s + 2k_f) - 2\phi(k_f - k_s)}{(k_s + 2k_f) + \phi(k_f - k_s)} \frac{1}{Pr} \frac{\partial^2 \theta}{\partial y^2}, \tag{8}$$

with the new boundary conditions as

$$\begin{aligned}
 u = 0, \quad v = 0, \quad \frac{\partial \theta}{\partial y} = -1, \quad \text{at} \quad y = 0, x \geq 0, \\
 u = u_e(x), \quad \frac{\partial u}{\partial y} = 0, \quad \theta = 0, \quad \text{as} \quad y \rightarrow \infty, x \geq 0,
 \end{aligned} \tag{9}$$

where $Pr = \mu_f C_p / k_f$ is Prandtl number, $M = \sigma B_0^2 a / \rho_f U_\infty$ is magnetic field, $K = k_o U_\infty / \mu_f a$ is viscoelastic parameter, and λ is mixed convection parameter.

3. Mathematical Solution

In order to solve Eqs. (6) to (8), subject to the boundary conditions (9), the following variables have been considered

$$\psi = xF(x, y), \quad \theta = \theta(x, y), \tag{10}$$

are introduced where ψ is the stream function defined as

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x}. \tag{11}$$

By substituting Eqs. (10) and (11) into Eqs. (6) to (8), obtained

$$\begin{aligned}
 & \left[(1-\phi) + \phi \frac{\rho_s}{\rho_f} \right] \left[\left(\frac{\partial F}{\partial y} \right)^2 + x \frac{\partial F}{\partial y} \left(\frac{\partial^2 F}{\partial x \partial y} \right) - x \frac{\partial F}{\partial x} \frac{\partial^2 F}{\partial y^2} - F \frac{\partial^2 F}{\partial y^2} \right] \\
 &= \left[(1-\phi) + \phi \frac{\rho_s}{\rho_f} \right] \frac{\sin x \cos x}{x} + \frac{1}{(1+\phi)^{2.5}} \frac{\partial^3 F}{\partial y^3} + \left[(1-\phi) + \phi \frac{(\rho\beta)_s}{(\rho\beta)_f} \right] \lambda \theta \frac{\sin x}{x} - M \frac{\partial F}{\partial y} \\
 &+ K \left[2 \frac{\partial F}{\partial y} \frac{\partial^3 F}{\partial y^3} - F \frac{\partial^4 F}{\partial y^4} - \left(\frac{\partial^2 F}{\partial y^2} \right)^2 + x \left(\frac{\partial^2 F}{\partial x \partial y} \frac{\partial^3 F}{\partial y^3} - \frac{\partial F}{\partial x} \frac{\partial^4 F}{\partial y^4} + \frac{\partial F}{\partial y} \frac{\partial^4 F}{\partial x \partial y^3} - \frac{\partial^2 F}{\partial y^2} \frac{\partial^3 F}{\partial x \partial y^2} \right) \right],
 \end{aligned} \tag{12}$$

$$\left[(1-\phi) + \phi \frac{(\rho C_p)_s}{(\rho C_p)_f} \right] \left[x \frac{\partial F}{\partial y} \frac{\partial \theta}{\partial x} - x \frac{\partial F}{\partial x} \frac{\partial \theta}{\partial y} - F \frac{\partial \theta}{\partial y} \right] = \frac{k_{nf}}{k_f} \frac{1}{Pr} \frac{\partial^2 \theta}{\partial y^2}, \quad (13)$$

which are subject to the following boundary conditions

$$F = 0, \quad \frac{\partial F}{\partial y} = 0, \quad \frac{\partial \theta}{\partial y} = -1, \quad \text{at } y = 0, x \geq 0, \quad (14)$$

$$\frac{\partial F}{\partial y} = \frac{\sin x}{x}, \quad \frac{\partial^2 F}{\partial y^2} = 0, \quad \theta = 0, \quad \text{as } y \rightarrow \infty, x \geq 0,$$

When $x \approx 0$, Eqs. (12) and (13) reduce to the following ordinary differential equations:

$$\frac{1}{(1+\phi)^{2.5}} f''' - \left[(1-\phi) + \phi \frac{\rho_s}{\rho_f} \right] [f'^2 - ff''] + K(2ff''' - ff^{iv} - f'^2) + \left[(1-\phi) + \phi \frac{(\rho\beta)_s}{(\rho\beta)_f} \right] \lambda \theta - Mf' = 0, \quad (15)$$

$$\frac{(k_s + 2k_f) - 2\phi(k_f - k_s)}{(k_s + 2k_f) + \phi(k_f - k_s)} \frac{1}{Pr} \theta'' + \left[(1-\phi) + \phi \frac{(\rho C_p)_s}{(\rho C_p)_f} \right] f \theta' = 0, \quad (16)$$

with the boundary conditions

$$f(0) = 0, \quad f'(0) = 0, \quad \theta'(0) = -1, \quad (17)$$

$$f'(\infty) = 1, \quad f''(\infty) = 0, \quad \theta(\infty) = 0,$$

4. Results

The behavior of fluid flow for viscoelastic nanofluid affected by magnetic field M is analyzed. The validation for numerical solutions is done by comparing the skin friction and heat transfer coefficients with the results from Nazar *et al.* [18]. Good agreements have been obtained from the results as shown in Table 2.

Table 2

Comparison values of heat transfer coefficient when $K = 0$, $Pr = 1$, $\phi = 0$, $M = 0$, and different values of λ

λ	Nazar <i>et al.</i> [18]	Present $f''(0)$	Nazar <i>et al.</i> [18]	Present $\theta(0)$
-0.2	1.0340	1.033028	1.8157	1.816890
0.4	1.5747	1.573759	1.7018	1.702823
3.0	2.4913	2.489892	1.4015	1.402232
10.0	5.7730	5.777805	1.1770	1.178456

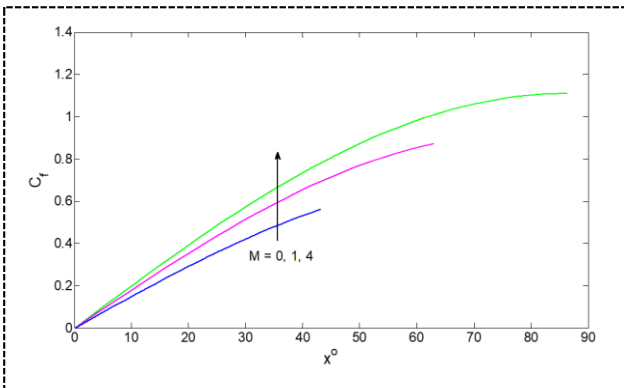


Fig. 1. Variation values of M for skin friction coefficient

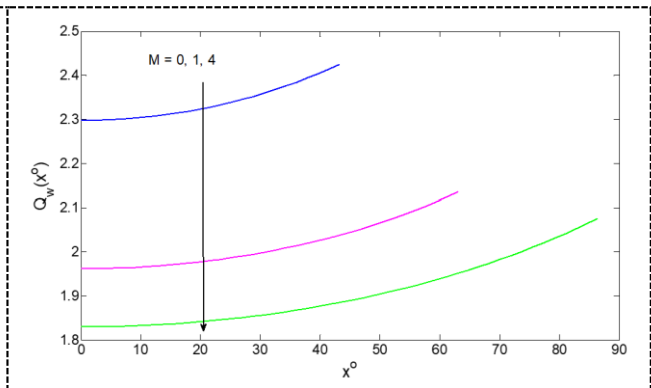


Fig. 2. Variation values of M for heat transfer coefficient

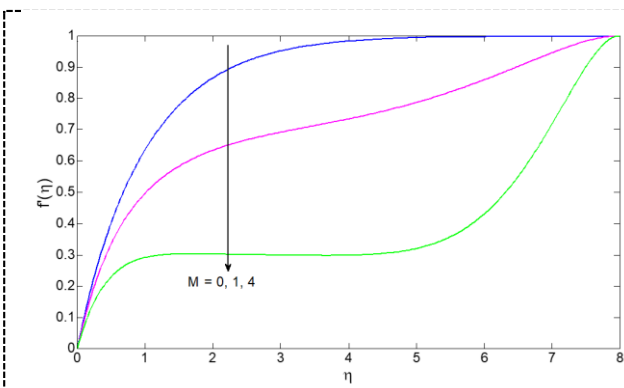


Fig. 3. Variation values of M for velocity profile

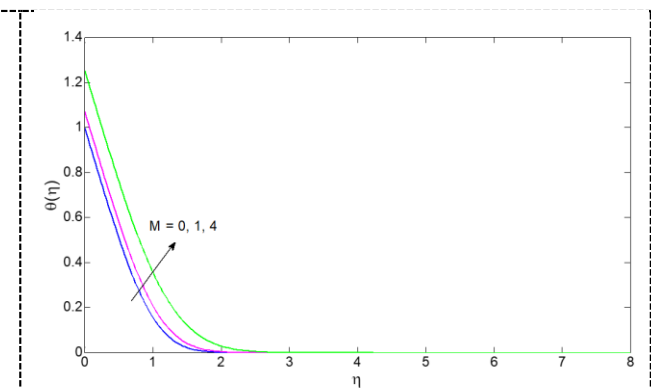


Fig. 4. Variation values of M for temperature profile

Figures 1–2 demonstrate how skin friction and heat transfer coefficients vary with M . Skin friction rises as M increases, as shown in Figure 1. Figure 2 shows, however, that heat transport reduces as M increases. Figures 3 and 4 depict the effects of M on velocity and temperature curves. As the value of M increases, the fluid velocity decreases (see Figure 3). Lorentz force is created by the existence of a transverse magnetic field that acts parallel to the flow. The flow is then slowed due to the resistance caused by this force. It's worth noting that this force has an unusual effect. It's worth noting that this force causes the fluid to suffer resistance by increasing friction between its layers, which causes the flow's temperature to rise, as seen in Figure 4. The temperature profile climbs significantly as the M value rises.

4. Conclusions

The conclusions drawn from this modern research are as follows: Velocity slows down and increase with augmentation in magnetic parameters. In the meantime, skin friction increases while heat transfer coefficient decreases for the increasing values of M .

Acknowledgement

The authors would like to acknowledge the Ministry of Higher Education (MOHE) for the financial support through vote numbers, FRGS/1/2019/STG06/UNIKL/03/1.

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