

Homann Stagnation Point Flow and Heat Transfer of Hybrid Nanofluids Over a Permeable Radially Stretching/Shrinking Sheet

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ABSTRACT

The steady two-dimensional Homann stagnation point flow and heat transfer of water base hybrid nanofluids (Al_2O_3 & Cu) over a permeable radially stretching/shrinking sheet have been studied. The similarity variables are introduced to transform the partial differential equations of the model into the ordinary differential equations. Numerical findings and dual solutions have been carried out by implementing the `bvp4c` code through MATLAB software. The most prominent effect is illustrated in the boundary layer thickness where the velocity profile increases upon the increment of the suction intensity but decreases in the temperature profile. Besides, the reduced Nusselt number also decreases as enlarging the value of copper and alumina nanoparticle volume fraction. The analysis of the first and second solutions are presented graphically with critical values as well as the detail discussions on the effects of the other involving parameters.

1. Introduction

The properties of heat transfer rate are always the attractive scope of research in fluid dynamics analysis. Since the conventional fluid has weak performance in the thermal conductivity, thus it has been alternatively replaced by the fluid that's more desirable and satisfactory, called nanofluid [1,2]. The discovery of nanofluid was coined by Choi and Eastman [3] who introduce it as an ideal solution to replace the conventional fluid for enhancing the thermal conductivity. Hence, intrigued by the numerous applications of nanofluids in diverse industries like power plant cooling system, micro-electronics, chemical engineering and so on, many researchers have work on the uniqueness of nanofluids [4-8].

However, lately the researchers have noticed that the nanofluids no longer to meet all the necessity in heat transfer properties. The inadequacy of conventional heat transfer nanofluids containing a single type of nanoparticle leads to the invention of hybrid nanofluid. As the extension

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work from nanofluids, the hybrid nanofluids are created with the amalgamation of distinct types of nanoparticles in the base fluid. The matches of different nanoparticles enable the great performance in thermal characteristics as well as ultimately increase the heat transfer rate. These were supported by many researchers after studied and investigated the attribute and behaviour of nanofluids in different aspects. Kumar and Arasu [9] claimed that composite nanoparticles exhibited superior improvement in rheological and thermal features of the conventional fluid compared to the single type nanoparticle after carried out the experimental research about stability, characterization and properties of $\text{Al}_2\text{O}_3\text{-Cu}$ hybrid nanofluids. Also, Hayat and Nadeem [10] claimed that the performance of Ag-CuO hybrid nanofluid in heat transfer is better than the mono-particle nanofluid. Hybrid nanofluids can be used in numerous conceivable applications in most of heat transfer fields like acoustics, manufacturing, micro fluidics etc. Due to the synergistic effect of choosing a proper combination of different nanoparticles and as the nanoparticles are being disseminate pertinently, it could lead to provide the exceptional high effective thermal conductivity. An extensive study on the related field had been done by a few researchers [11-13]. Sinz *et al.*, [14] discovered an increment of heat transfer rate in the turbulent force convective model of silver-graphene hybrid nanofluid as Reynolds number increases. The comprehensive review on the preparation methods of hybrid nanofluid with its properties and performance had been carried out by Xian *et al.*, [15] and Urmi *et al.*, [16]. Azman *et al.*, [17] compared the heat transfer performance between the mono nanofluid (single nanoparticle) and hybrid nanofluid (two nanoparticles, $\text{Al}_2\text{O}_3\text{-Cu}$) in a straight pipe. They found that the percentage of mono nanofluid heat transfer was slightly higher than the hybrid nanofluid. In contrast, Idris *et al.*, [18] discovered that the hybrid nanofluid ($\text{Al}_2\text{O}_3\text{-SiO}_2$) improved the single nanofluid performance in the cooling plate of fuel cells.

The heat transfer and fluid flow due to stretching plate is one of the crucial parts in the applications like metal or fibre spinning manufacturing, paper production, plastic films, glass blowing and etc. The amount of heat transfer at the stretching surfaces is directly affecting the quality of the final products as mentioned by Roşca and Pop [19]. The first problem which involved the boundary layer flow on a continuous object's surface was studied by Sakiadis [20], followed by Crane [21] who studied the flow past a stretching plate. He came out with the exact analytical solution for the steady two-dimensional flow over a stretching surface. Thus, this has inspired many researchers to consider stretching sheet or plate in their boundary layer study. The stretching flow can be identified as a circumstance when the fluid at the boundary moves away from a fixed point whereas if the movement of the fluid against the direction of the flow, then it is called shrinking flow. Miklavčič and Wang [22] proved the existence of non-uniqueness solution of the fluid flow owing to a shrinking sheet. After several years, most of the researchers considered both stretching/shrinking sheet in their studies. The stagnation-point flow over a stretching/shrinking sheet in a nanofluid was analysed by Bachok *et al.*, [23]. Then, Waini *et al.*, [24] investigated the unsteady flow and heat transfer of hybrid nanofluid past a stretching/shrinking sheet. The solutions for greater shrinking impermeable sheet will never exist since the vorticity could not be cramped in the boundary layer. However, if there is the stagnation flow to restrain the vorticity, thus the similarity solutions are likely to exist.

The problem which involved the two-dimensional stagnation-point flows was pioneered by Hiemenz [25]. The flow of stagnation point can be defined as the flow about the front of a blunt-nosed body or the stagnation division, happen on bodies moving in a fluid. Moreover, the stagnation point meets unexpectedly the highest rates of mass deposition, the topmost heat transfer as well as the pressure. Then, Homann [26] studied the problem of stagnation-point flow in axisymmetric case. Afterwards, the study of numerous aspects of stagnation-point flow in large scale garnered considerable attention which correspond to the viscous fluid flow and different geometry like stretching/shrinking plate and wavy cylinder. Yousefi *et al.*, [27] investigated the steady general

three-dimensional stagnation-point flow of an aqueous $\text{TiO}_2\text{-Cu}$ hybrid nanofluid toward a wavy cylinder analytically. Besides, Ghalebaz *et al.*, [28] studied the mixed convection and stagnation point flow and heat transfer of $\text{Al}_2\text{O}_3\text{-Cu}$ /water hybrid nanofluid over a vertical plate. The simultaneous effects of wall transpiration through a radially stretching on Homann stagnation point flow was reported by Weidman and Ma [29]. They mentioned that the blowing must be less than zero while suction must be greater than zero. Recently, Mahapatra and Sidui [30] investigated the non-axisymmetric Homann stagnation-point flow of a viscoelastic fluid past a rigid plate and concluded that the fluid's viscoelasticity delayed the reverse flow at the wall.

Motivated by the above researches, the present studies intend to theoretically carry out numerical solutions of Homann stagnation point flow and heat transfer of $\text{Al}_2\text{O}_3\text{-Cu}$ hybrid nanofluids past a permeable radially stretching/shrinking sheet. The novelty of our work is focus on the dual solutions of Homann stagnation point flow of hybrid nanofluid. This topic yet to fully explored and discovered and it is important to imply in the specific field as mentioned above. Since some of the researchers only provide a single solution for particular problem, thus may have resulted in a lack of persuasion and cause the explanation looks insubstantial. By performing the dual solutions, we are able to see the differences of the two obtained results and enable to give adequate explanation for conviction about the behaviour of hybrid nanofluids flow. The graphical results are presented and plotted, then discussed in the next sections.

2. Mathematical Model

In this study, the model of steady two-dimensional Homann stagnation point flow and heat transfer of the water base hybrid nanofluid over a permeable stretching/shrinking sheet is considered as displayed in Figure 1 with the cylindrical coordinates (r, Ω, z) . Here, the plate is reflected by the r -axis and z axis is perpendicular to it. All the physical quantities are not depending on the angle Ω since it is rotational symmetry. It is assumed that $u_w(r)$ is representing the velocity of the stretching/shrinking sheet. In addition, $(u_e(r), 0, 0)$ is the velocity of the far field (inviscid flow). As further information, the constant temperature of the surface of the sheet is denoted by T_w while T_∞ represents the temperature of the far field. Since this study focus on hybrid nanofluid, thus a few assumptions need to be done which are considering the uniformity size of nanoparticles and neglecting the impact of the nanoparticle's agglomeration. This is due to the synthetization of the nanofluids as a stable mixture of the base fluid and nanoparticles. Therefore, with these physical assumptions, we consider the cylindrical coordinates $(r, 0, z)$ of the governing equations of the Homann stagnation flow of hybrid nanofluid as follows [29,31]:

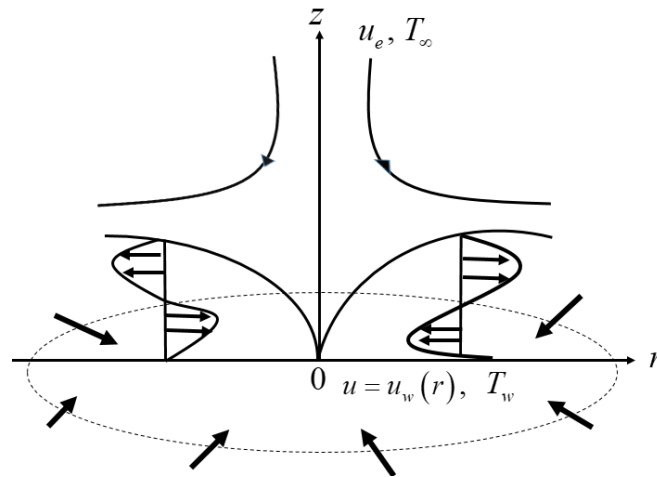


Fig. 1. Physical model of the Homann stagnation point flow of hybrid nanofluid over a radially shrinking sheet ($\lambda < 0$)

$$\frac{\partial}{\partial r}(ru) + \frac{\partial}{\partial z}(rw) = 0 \quad (1)$$

$$u \frac{\partial u}{\partial r} + w \frac{\partial u}{\partial z} = -\frac{1}{\rho_{hmf}} \frac{\partial p}{\partial r} + \frac{\mu_{hmf}}{\rho_{hmf}} \left(\frac{\partial^2 w}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{\partial^2 u}{\partial z^2} - \frac{u}{r^2} \right) \quad (2)$$

$$u \frac{\partial T}{\partial r} + w \frac{\partial T}{\partial z} = \frac{k_{hmf}}{(\rho C_p)_{hmf}} \left(\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \frac{\partial T}{\partial r} + \frac{\partial^2 T}{\partial z^2} \right) \quad (3)$$

subject to

$$\begin{aligned} w = w_0, \quad u = u_w(r) = br, \quad T = T_w \quad \text{at} \quad z = 0 \\ u = u_e(r) \rightarrow ar, \quad T \rightarrow T_\infty \quad \text{as} \quad z \rightarrow \infty. \end{aligned} \quad (4)$$

Here $(u, 0, w)$ represent the velocity components of the hybrid nanofluid along $(r, 0, z)$ axes, w_0 denotes the constant mass flow, T is the temperature of the hybrid nanofluid and a is the positive strain rate of the stagnation-point flow. Further, μ_{hmf} is the dynamic viscosity of the hybrid nanofluids, ρ_{hmf} is the density of the hybrid nanofluids, k_{hmf} is the thermal conductivity of the hybrid nanofluid and $(\rho C_p)_{hmf}$ is the heat capacity of the hybrid nanofluid. The physical features of nanofluids and hybrid nanofluids are given in Table 1 where ϕ is the nanoparticle volume fraction ($\phi = 0$ reflects the base fluid), ρ_f and ρ_s are the densities of the base fluid and the hybrid nanoparticle, respectively, C_p is the heat capacity at the constant pressure, $(\rho C_p)_f$ and $(\rho C_p)_s$ are the heat capacitance of the base fluid and the hybrid nanoparticle, respectively whereas k_f and k_s are the thermal conductivities of the base fluid and the hybrid nanoparticles, respectively. Then, the physical properties of copper Cu and alumina Al_2O_3 nanoparticles as well as the base fluid (water) are provided in Table 2.

Table 1
 Physical features of conventional nanofluids and hybrid nanofluids

Properties	Nanofluid	Hybrid nanofluid
Density	$\rho_{nf} = (1-\phi)\rho_f + \phi\rho_s$	$\rho_{hnf} = \{(1-\phi_2)[(1-\phi_1)\rho_f + \phi_1\rho_{s1}]\} + \phi_2\rho_{s2}$
Heat Capacity	$(\rho C_p)_{nf} = (1-\phi)(C_p)_f + \phi(\rho C_p)_s$	$(\rho C_p)_{hnf} = \{(1-\phi_2)[(1-\phi_1)(\rho C_p)_f + \phi_1(\rho C_p)_{s1}]\} + \phi_2(\rho C_p)_{s2}$
Viscosity	$\mu_{nf} = \frac{\mu_f}{(1-\phi)^{2.5}}$	$\mu_{hnf} = \frac{\mu_f}{(1-\phi_1)^{2.5} \cdot (1-\phi_2)^{2.5}}$
Thermal Conductivity	$k_{nf} = \frac{k_{s1} + 2k_f - 2\phi_1(k_f - k_{s1})}{k_{s1} + 2k_f + \phi_1(k_f - k_{s1})} \times k_f$	$k_{hnf} = \frac{k_{s2} + 2k_{nf} - 2\phi_2(k_{nf} - k_{s2})}{k_{s2} + 2k_{nf} + \phi_2(k_{nf} - k_{s2})} \times k_{nf}$ where $k_{nf} = \frac{k_{s1} + 2k_f - 2\phi_1(k_f - k_{s1})}{k_{s1} + 2k_f + \phi_1(k_f - k_{s1})} \times k_f$

Table 2
 Thermophysical features of the nanoparticles and base fluid

Properties	Water (f)	Al ₂ O ₃ (ϕ_1)	Cu (ϕ_2)
$\rho(kg / m^3)$	997.0	3970	8933
$C_p(J / kgK)$	4180	765	385
$k(W / mK)$	0.6071	40	400

We introduce now, the following similarity variables [29]

$$u = arf'(\eta), \quad w = -2\sqrt{av_f}f(\eta), \quad \theta(\eta) = \frac{T - T_\infty}{T_w - T_\infty}, \quad \eta = z\sqrt{\frac{a}{v_f}} \tag{5}$$

where prime symbolizes differentiation with respect to η . Furthermore,

$$w_0 = -\sqrt{av_f}S \tag{6}$$

where S is the mass flux parameter with for $S > 0$ suction and $S < 0$ for injection, respectively. By implying the similarity variables (5), Eq. (2) and Eq. (3) are reduced into the ordinary differential equations as follows

$$\frac{\mu_{hnf}/\mu_f}{\rho_{hnf}/\rho_f} f''' + 2ff'' - f'^2 + 1 = 0 \tag{7}$$

$$\frac{1}{Pr} \frac{k_{hnf}/k_f}{(\rho C_p)_{hnf}/(\rho C_p)_f} \theta'' + 2f\theta' = 0 \tag{8}$$

correspond to the boundary conditions

$$\begin{aligned} f(0) = S, \quad f'(0) = \lambda, \quad \theta(0) = 1 \\ f'(\eta) \rightarrow 1, \quad \theta(\eta) \rightarrow 0 \quad \text{as} \quad \eta \rightarrow \infty \end{aligned} \quad (9)$$

where $\text{Pr} = \frac{\mu_f (C_p)_f}{k_f}$ Prandtl number. Besides, $\lambda = \frac{b}{a}$ is the stretching/shrinking parameter where $\lambda > 0$ represents stretching sheet while $\lambda < 0$ physically means the sheet is shrunk. In addition, $\lambda = 0$ shows that the sheet is static.

The physical quantities of interest are the radial skin friction along the surface of the sheet $\bar{\tau}_r$ and the heat flux from the surface of the sheet \bar{q}_r , which are defined as

$$\bar{\tau}_r = \mu_{hmf} \left(\frac{\partial u}{\partial z} \right)_{z=0}, \quad \bar{q}_r = -k_{hmf} \left(\frac{\partial T}{\partial z} \right)_{z=0} \quad (10)$$

Using Eq. (5), we get the dimensionless radial skin friction along the sheet τ_r and dimensionless heat flux from the sheet, as

$$\tau_r = \frac{\mu_{hmf}}{\mu_f} f''(0), \quad q_r = -\frac{k_{hmf}}{k_f} \theta'(0) \quad (11)$$

$$\text{where } \tau_r = \frac{\bar{\tau}_r}{\rho_f \mu_f \sqrt{v_f a^3} r} \quad \text{and} \quad q_r = \frac{\bar{q}_r}{k_f (T_w - T_\infty) \sqrt{a/v_f}}.$$

3. Results and Discussions

As the first step, it is necessary to transform the partial differential Eq. (1)-(3) into the first order system of ordinary differential equations. After that, the ordinary differential equations for Eq. (7) and Eq. (8) together with boundary conditions (9) were numerically solved by executing the programming codes of boundary value problem solver (bvp4c) through MATLAB software in order to obtain the numerical solutions. The exact dual solutions only been obtained until the asymptotic results appeared through selection of an appropriate initial guess value. Then, the graphical results for the velocity profile, temperature profile, reduced skin friction coefficient and reduced Nusselt number for each involved parameters such as suction parameter and nanoparticle volume fraction (ϕ_1 and ϕ_2) are generated. Figure 2 and Figure 3 portray the influence of suction parameter on the velocity and temperature profiles respectively. It is noticed that Figure 2 shows a rising of the velocity profile of hybrid nanofluids as an escalation in the value of suction parameter but the second solution shows the opposite behaviour which reducing the fluid flow's velocity profile. In Figure 3, both solutions are showing the same trend where the thermal boundary layer thickness is decreasing when the greater suction parameter imposed.

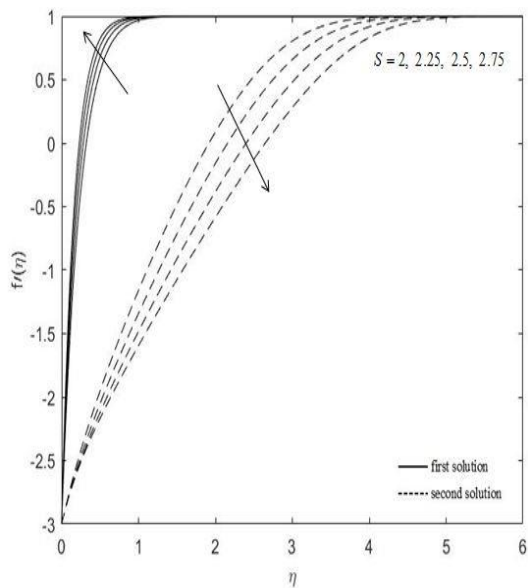


Fig. 2. The trend of velocity profile $f'(\eta)$ with variation in S

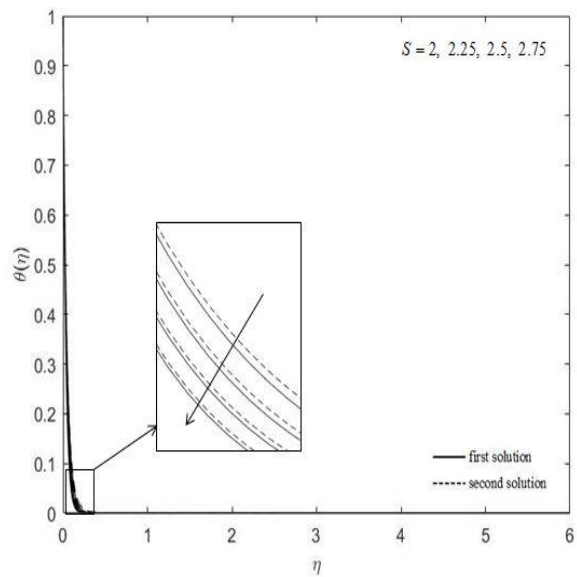


Fig. 3. The trend of temperature profile $\theta(\eta)$ with variation in S

The trend of velocity profile with variation in Cu nanoparticle volume fraction decreases for second solution but increase for the first solution as presented in Figure 4. Due to the rise in shear stress and subsequently shrinks the momentum boundary layer thickness, the fluid flow velocity is declining. The impact of variations in Cu nanoparticles concentration on the temperature profile showing the prominent increment for both solutions as exhibited in Figure 5. Since the convective heat transfer has been reduced, the heat is then transfer from heated hybrid nanofluid to the hot uppermost layer and resulted temperature increase in the both solutions.

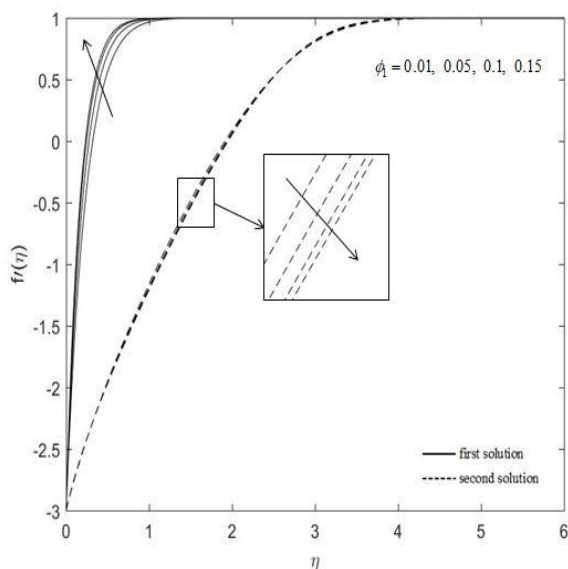


Fig. 4. The trend of $f'(\eta)$ with variation in Cu nanoparticle volume fraction

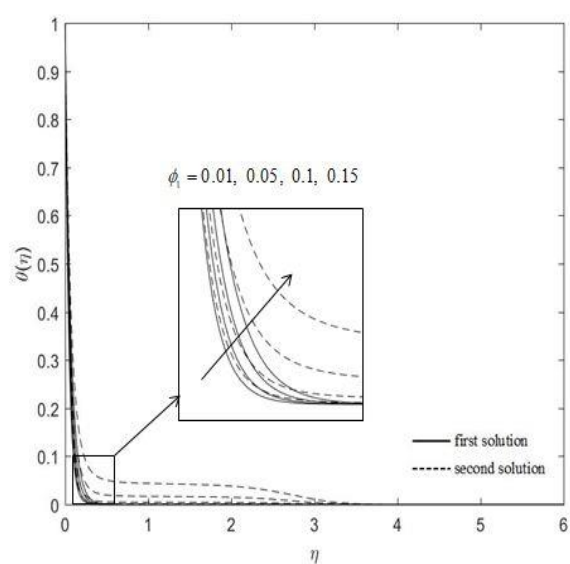


Fig. 5. The trend of $\theta(\eta)$ with variation in Cu nanoparticle volume fraction

On the other hand, Figure 6 and Figure 7 illustrate the trends of velocity and temperature profile for alumina nanoparticle volume fraction, respectively. The velocity profile increases in second solutions but having opposite behaviour in the first solution as appeared asymptotically in Figure 6 and it is noticed that the thickness of momentum boundary layer for the first solution is lower if compared to the second solution. Incidentally, the profile of temperature displays the same trend in which increases as the alumina nanoparticle volume fraction increases. It reveals that the thickness of thermal boundary layer is getting higher with the increase of alumina nanoparticles.

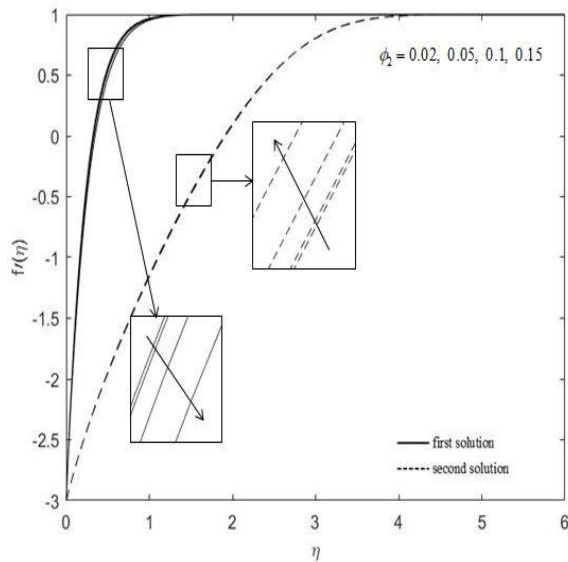


Fig. 6. The trend of $f'(\eta)$ with variation in Al_2O_3 nanoparticle volume fraction

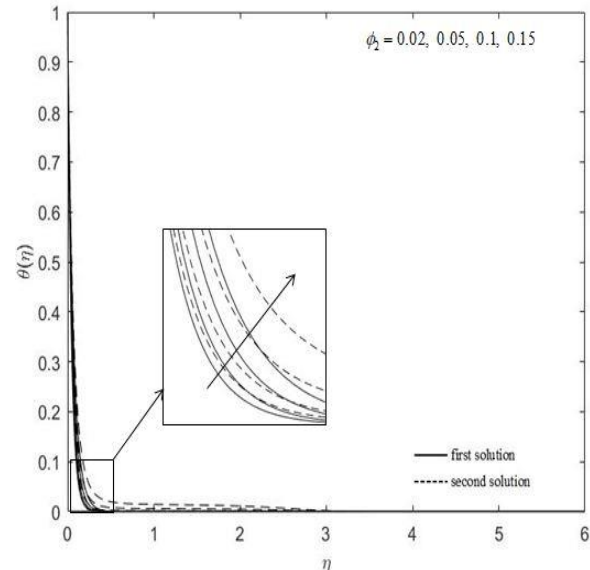


Fig. 7. The trend of $\theta(\eta)$ with variation in Al_2O_3 nanoparticle volume fraction

Figure 8 and Figure 9 depict the influence of suction parameter on reduced skin friction coefficient and reduced Nusselt number for dual solutions, respectively. As shown in the both figures, we noticed that there is a unique solution exists up for different values of suction parameter. Furthermore, there is no existence of the solutions if the λ less than λ_c . The values of λ_c are -4.3259, -5.8191 and -7.5934 which correspond to S equals to 2, 2.5 and 3, respectively. It is apparently shown that the increment of S which leads to rise of the boundary layer separation in dual solutions, increasing $|f''(0)|$ at all the time as in Figure 8. Since the suction gives rise to more resistance in the fluid flow and causes the thickness of boundary layer become thicken, thus consequently the skin friction coefficient increases rapidly. However, the first solution shows us that the skin friction coefficient drops back slightly and retard once getting closer with the λ_c . There is a peculiar behaviour shown in Figure 8 (second solution for $S=2.5$) and this is due to the retardation at specific values of λ . For Figure 9, the reduced Nusselt number decrease slowly as increment of suction parameter. The temperature gradient decreases and reduces the thickness of boundary layer which help the hybrid nanofluid to transfer the heat effortlessly. Also, it is worth to notice that the $-\theta'(0)$ has slow response in the beginning but increases extremely starting from $\lambda < -2$.

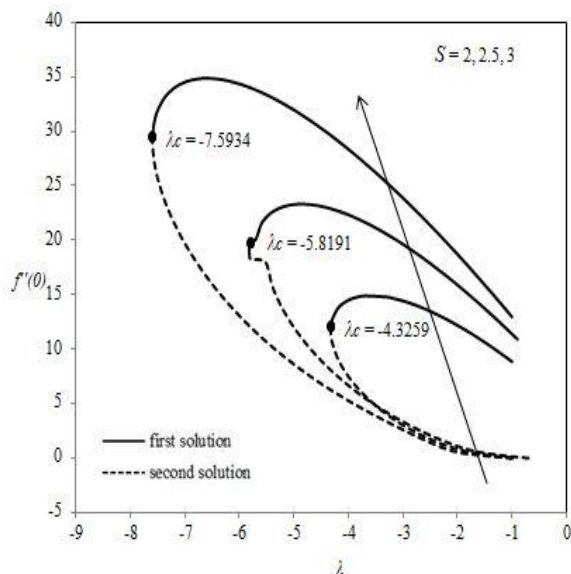


Fig. 8. The trend of $f''(0)$ with λ for several values of S

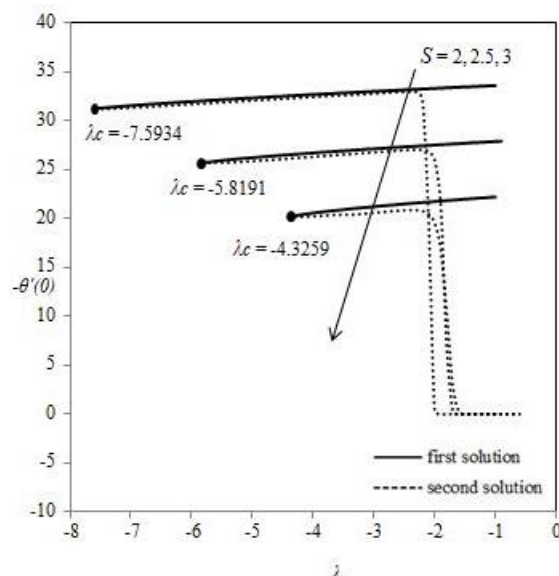


Fig. 9. The trend of $-\theta'(0)$ with λ for various value of S

Figure 10 and Figure 11 illustrate the effect of copper nanoparticle volume fraction ϕ_1 on $|f''(0)|$ and $-\theta'(0)$ respectively. The values of λ_c are -5.5369, -5.9501 and -6.6125 which correspond to ϕ_1 equals to 0.05, 0.1 and 0.15, respectively. In Figure 10, the greater value of ϕ_1 reflects to the lower value of λ_c . In addition, by adding up the value of ϕ_1 results in the increase of $|f''(0)|$ for all the time in both solutions. Normally, the increment of nanoparticle volume fraction indicates a rapid increase of the movement and vibration of the nanoparticles in the fluid flow, thus expecting more heat energy to be released. Nevertheless, this phenomenon is contradictory proven in Figure 11 where the reduced Nusselt number lessen as ϕ_1 getting higher. The heat transfer rate decreases slowly when the nanoparticles concentration increases.

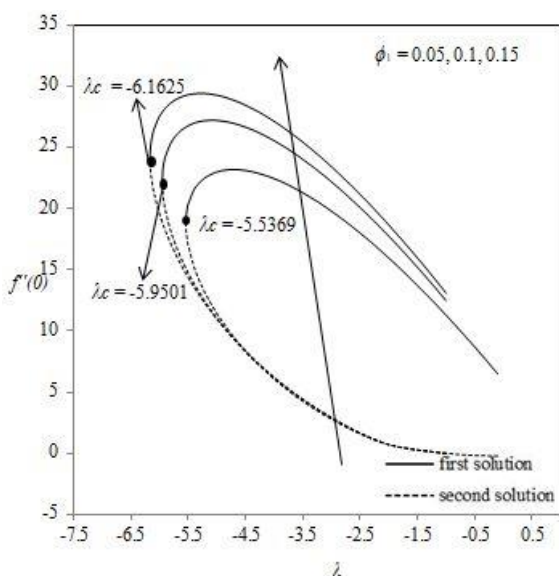


Fig. 10. Variations of $f''(0)$ for several value of Cu nanoparticle volume fraction

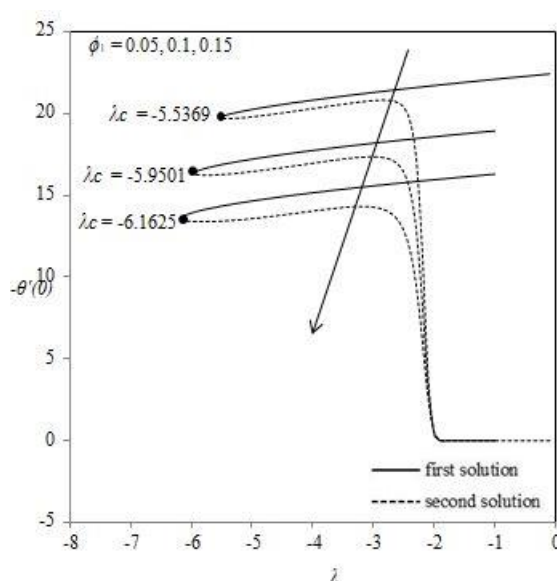


Fig. 11. Variation of $-\theta'(0)$ for several value of Cu nanoparticle volume fraction

Figure 12 and Figure 13 present the effect of alumina nanoparticle volume fraction ϕ_2 on $|f''(0)|$ and $-\theta'(0)$ respectively. The values of λ_c are -5.0245, -4.9476 and -4.8120 where correspond to ϕ_2 equal to 0.05, 0.1 and 0.15, respectively. For Figure 12, it is noticed that the greater the value of ϕ_2 , the higher the value of λ_c which is opposite behaviour from copper. This means that the higher concentration of alumina contributes to the earlier boundary layer separation. Also, a rise in in ϕ_2 causes the decrement of $|f''(0)|$ for the dual solutions and slightly decreases the reduced Nusselt number as shown in Figure 13.

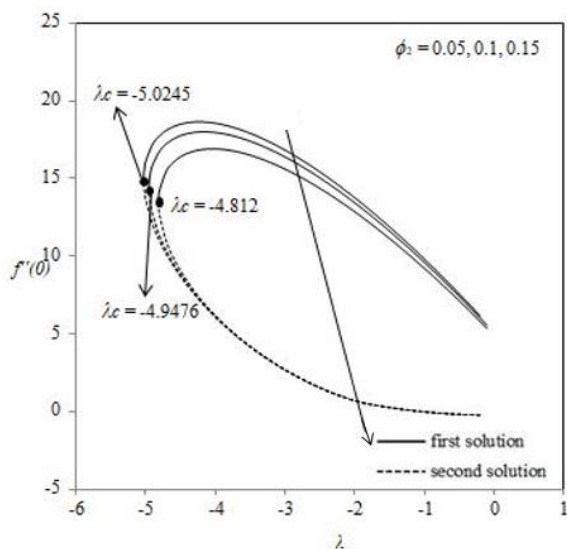


Fig. 12. Variations of $f''(0)$ for several value of Cu nanoparticle volume fraction

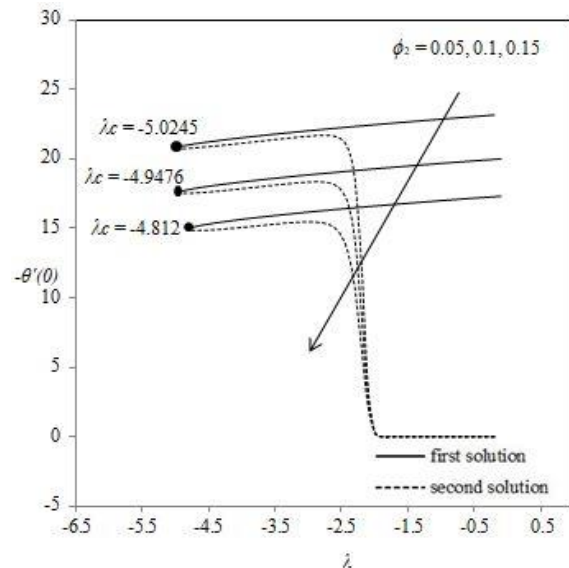


Fig. 13. Variation of $-\theta'(0)$ for several value of Cu nanoparticle volume fraction

4. Conclusions

The problem of steady two-dimensional Homann stagnation point flow and heat transfer of water-based hybrid nanofluids over a permeable radially stretching/shrinking sheet has been solved numerically. The important of finding results and prominent effects are summarised as follows

- i. The suction parameter caused the increment of velocity profile but in contrary reduced the temperature profile and thermal boundary layer thickness.
- ii. The velocity profiles for Cu and Al_2O_3 nanoparticle volume fraction showed the opposite trend in dual solutions.
- iii. Cu and Al_2O_3 nanoparticle volume fraction enlarged the thickness of boundary layer in temperature profile for both solutions.
- iv. The presence of higher concentration Cu nanoparticles led to a rise in the reduced skin friction coefficient but a contradict trend was observed for Al_2O_3 nanoparticles. However, both nanoparticles contributed to a decrement in the reduced Nusselt number.

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